Mapping Transition-edge Sensor Temperature Sensitivity and Current Sensitivity Surface as a Function of Current, Magnetic Field and Temperature with IV curve and Complex Impedance Measurement

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Motivation

1. To resolve the spectra of the astronomical diffuse X-ray background in the 0.1-0.5 keV energy range, we need large area pixels $(\sim 1 \text{mm}^2)$ with excellent energy resolution (1-2) eV).

<u>TES devices – overview^[5]</u>			TES behaves like a Josephson Junction (Shapiro Steps)[7]	
 Mo-Au bilayer No membrane/absorber 	³⁵ MXW0203 @ <i>I</i> _{bias} = 120uA 106.6mK 25 20 15 15 10	These two Mo-Au bilayer TESs are diagnosis chips with no membrane or absorber design. MXW0203 Tc ~ 108mK; MXW1213 Tc ~ 104mK. Here we take	A is a in the second se	XW0203 Jrve tak 106.671 n extern ternatin field is pplied to equenc 254.2961

2. For TES pixels with large absorbers that have a relatively high heat capacity, high α and low β are needed to achieve the required energy resolution.

3. Most work in the field has found high α to be correlated with high β and also with excess noise. The cause of this correlation remains to be understood[6].





















1. McCammon, Dan. "Thermal equilibrium calorimeters-an introduction." Cryogenic particle

detection (2005): 1-34.

2. Maasilta, Ilari J. "Complex impedance, responsivity and noise of transition-edge sensors: Analytical solutions for two-and three-block thermal models." Aip Advances 2.4 (2012): 042110.

3. Lindeman, M. A., et al. "Accurate thermal conductance and impedance measurements of transition edge sensors." Journal of Low Temperature Physics 151.1-2 (2008): 180-184.

4. Zhang, Shuo, et al. "Mapping of the resistance of a superconducting transition edge sensor as a function of temperature, current, and applied magnetic field." Journal of Applied Physics 121.7 (2017): 074503.

5. K. M. Morgan, S. E. Busch, M. E. Eckart, C. A. Kilbourne, D. McCammon, "Large Area Transition Edge Sensor X-ray Microcalorimeters for Diffuse X-ray Background Studies", JLTP 176, 331 (2013)

6. Jethava, Nikhil, et al. "Dependence of excess noise on the partial derivatives of resistance in superconducting transition edge sensors." AIP Conference Proceedings. Vol. 1185. No. 1. AIP, 2009.

7. J. Sadleir, Applied Superconductivity Conference, 2016

8. S. Shapiro, "Josephson currents in superconducting tunneling: The effect of microwaves and other observations" Physical Review Letters 11, 80 (1963). C. C. Grimes and S. Shapiro, Physical Review 169, 397 (1968).

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Summary

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1. We characterized two TES devices and compared them. The TES with no finger or bank structures is more sensitive to the magnetic field than the three finger TES, while the period of the fringes on the BI curve is about the same for both devices.

2. The IV curves of both devices show Shapiro steps when a high frequency external field is applied. More Shapiro steps are visible on the zero finger TES IV curve than the three finger TES.

3. The simple thermal model with only one time constant can't fit the complex admittance very well. The fitting results give smaller α compared with IV curve. The two block thermal model yields a better fit.